

Non-continuous interval maps and rotation theory in spiking neuron models

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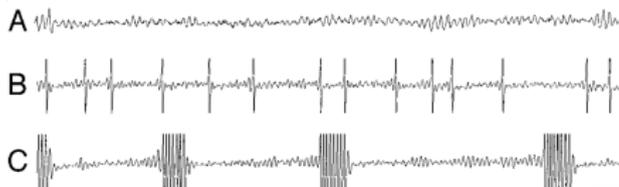
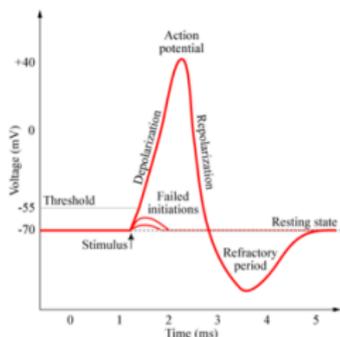


Będlewo; 28 May 2015

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- 3 Rotation theory
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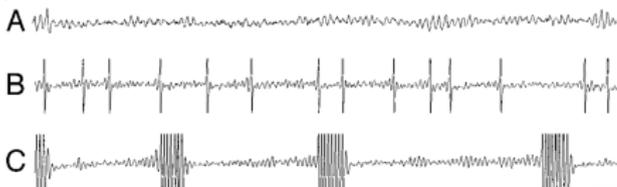
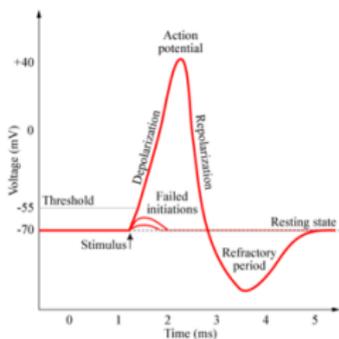
What characterizes neurons' activity? Neurons are **electrically excitable cells** that communicate through the emission of **action potentials (spikes)**: **stereotyped membrane potential electrical impulses**. They encode information in the way spikes are emitted through:

- the answer to specific simple stimuli (**excitability properties**)
- and the spike pattern fired,
- often related to properties of the interspike behavior (e.g.: sub-threshold oscillations) .



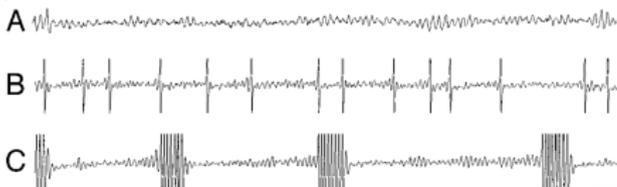
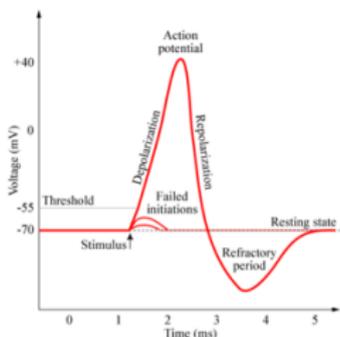
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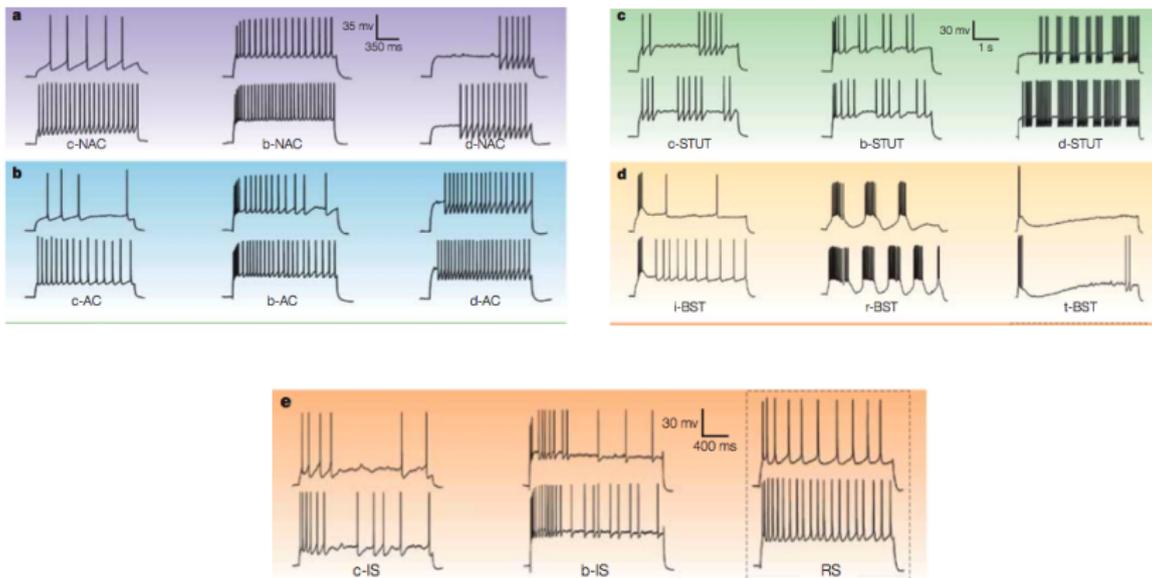


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Different spike patterns



Markram *et al* 2004

Phenomenological Neuron Models

Aimed to reproduce the typical behaviors of nerve cells in response to different stimuli:

- Excitability properties of neurons
- Frequency preference property
- subthreshold oscillations,
- Spike patterns fired

Neuron and Dynamical Systems

The main excitability properties can be linked with bifurcations of dynamical systems for

- **Continuous dynamical systems**: detailed neuron models and their reductions (Rinzel, Ermentrout, Guckenheimer, ...).
- **Discrete dynamical systems**: map-based models (Caselles, Rulkov, ...)

Hybrid dynamical systems

Integrate-and-fire neuron models combine:

- A **continuous** dynamical system (ordinary differential equations) accounting for input integration
- A **discrete** dynamical system (map iteration) accounting for spike emission.

Classical Integrate-and-Fire Neurons

$$\frac{dv}{dt} = -v + I$$

$$v = \theta \Rightarrow \text{Spike!}$$

Louis Lapicque, 1907

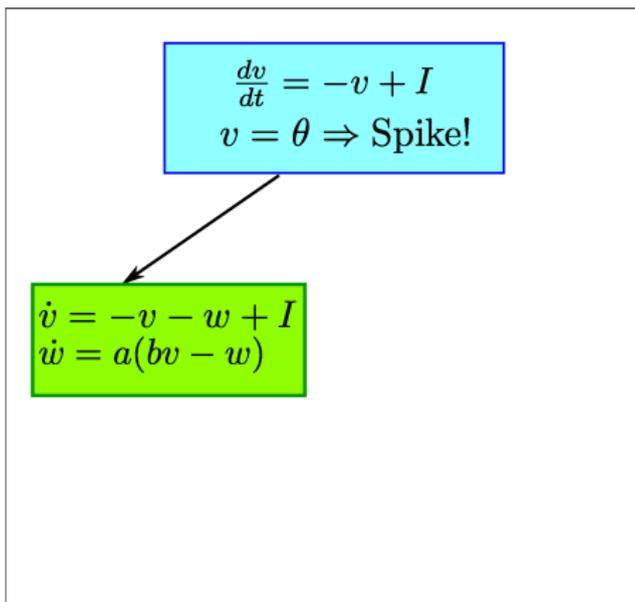
Izhikevich (2003)

$$\begin{cases} \dot{v} = v^2 - w + I \\ \dot{w} = a(bv - w) \end{cases}$$

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Classical Integrate-and-Fire Neurons



Wehmeier *et al*, 1989

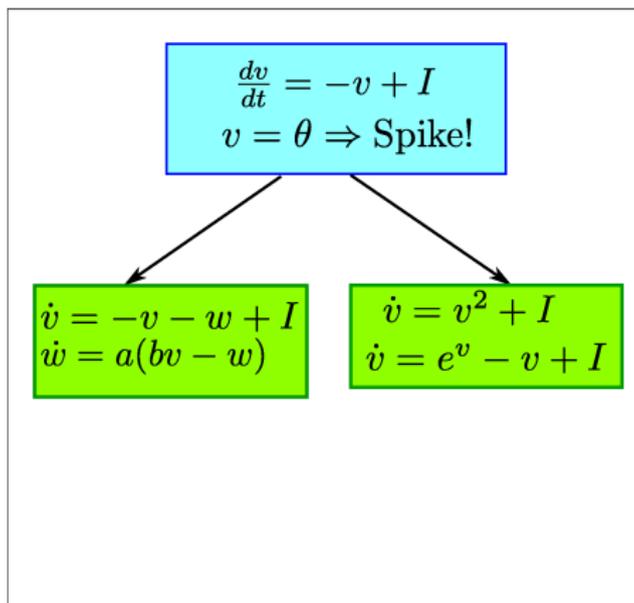
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e.g. Ermentrout Kopell, 1982, Fourcaud-Trocme *et al* 2003

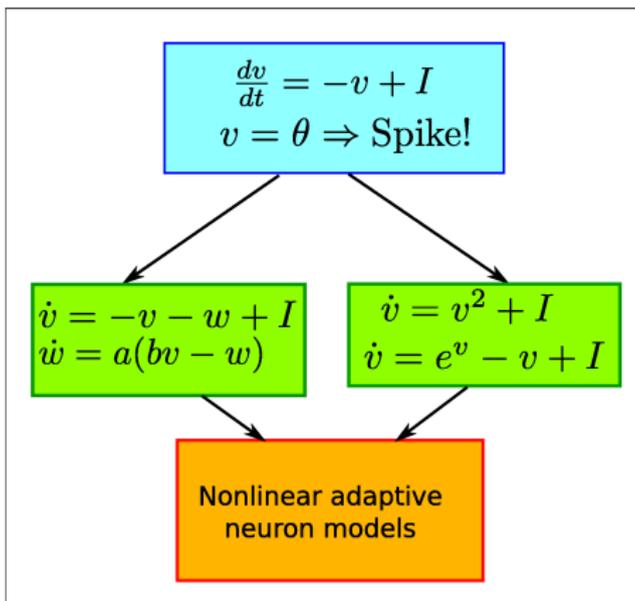
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Izhikevich, 2003 & Brette Gerstner 2005

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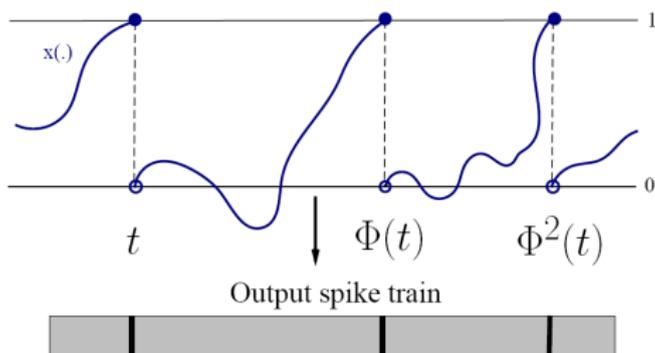
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One-dimensional models

$$\dot{x} = F(t, x) \quad F: \mathbb{R}^2 \rightarrow \mathbb{R}$$

$$\lim_{s \rightarrow t^+} x(s) = 0, \quad \text{if } x(t) = 1$$



Definition [Firing map]

$$\Phi: \mathbb{R} \rightarrow \mathbb{R}, \quad \Phi(t) = \min\{s > t : x(s; t, 0) = 1\}$$

$$D_\Phi = \{t \in \mathbb{R} : \exists_{s > t} x(s; t, 0) = 1\}$$

$$t_n = \Phi^n(t) = \min\{s > \Phi^{n-1}(t) : x(s; \Phi^{n-1}(t), 0) = 1\}$$

Perfect Integrator Model

$$\dot{x} = f(t) \quad (\text{PI})$$

Leaky Integrate-and-Fire

$$\dot{x} = -\sigma x + f(t) \quad (\text{LIF})$$

Non-linear models

$$\dot{x} = F(t, x)$$

- Mathematical analysis of one-dimensional IF models was performed e.g. in [R.Brette, 2004], [H.Carrillo, F.A.Ongay, 2001], [T.Gedeon, M.Holzer, 2004], [W. Marzantowicz, J.S., 2011], [W. Marzantowicz, J.S., 2015] and [P. Kasprzak, A. Nawrocki, J. S., 2015] (with focus on periodic and almost-periodic input functions)

Bidimensional integrate-and-fire models

Bidimensional integrate-and-fire models:

$$\dot{v} = F(v) - w + I \quad (1)$$

$$\dot{w} = a(bv - w) \quad (2)$$

A spike is emitted at time t^* such that

$$\lim_{t \rightarrow t^{*-}} v(t) = \infty$$

At the moment of the spike we reset:

$$v(t^{*+}) \rightarrow v_R, \quad w(t^{*+}) \rightarrow \gamma w(t^{*-}) + d$$

The adaptation map: $\Phi(w_0) = \gamma w(t^{*-}) + d$, (v_R, w_0) -the initial condition of the solution $(v(t), w(t))$ which spikes at t^* .

Examples include adaptive exponential model ($F(v) = -v + \gamma \exp(\beta v)$), generalized integrate-and-fire model ($F(v) = v^2$) and leaky integrate-and-fire model ($F(v) = v^2 + \lambda v$).

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Examples include *adaptive exponential model* ($F(v) = e^v - v$), *quadratic adaptive model* ($F(v) = v^2$) and *quartic model* ($F(v) = v^4 + 2av$)

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The model can display complex dynamics including **Mixed-Mode Oscillations** and **Mixed-Mode Bursting Oscillations** (MM(B)O) that are sequences of spikes interspersed by small subthreshold oscillations.

MM(B)Os so far have been investigated in 3D and higher dimensional systems ([M. Desroches et al., 2012], [M. Krupa et al., 2012], [T. Vo et al., 2012]).

In such hybrid models they have never been observed before.

From the neuroscience point of view, they have been evidenced in Hodgkin-Huxley model ([J. Rubin, M. Wechselberger, 2007], [J. Rubin, M. Wechselberger, 2008]) and in the coupled FitzHugh-Nagumo systems ([N. Berglund, D. Landon, 2012], [M. Desroches et al., 2008]).

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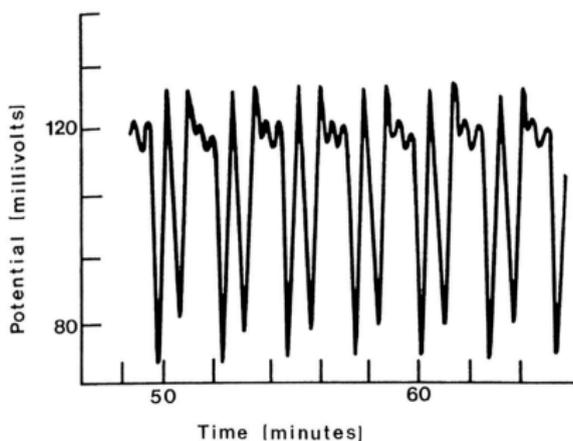
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Another (classical) example are chemical reactions:



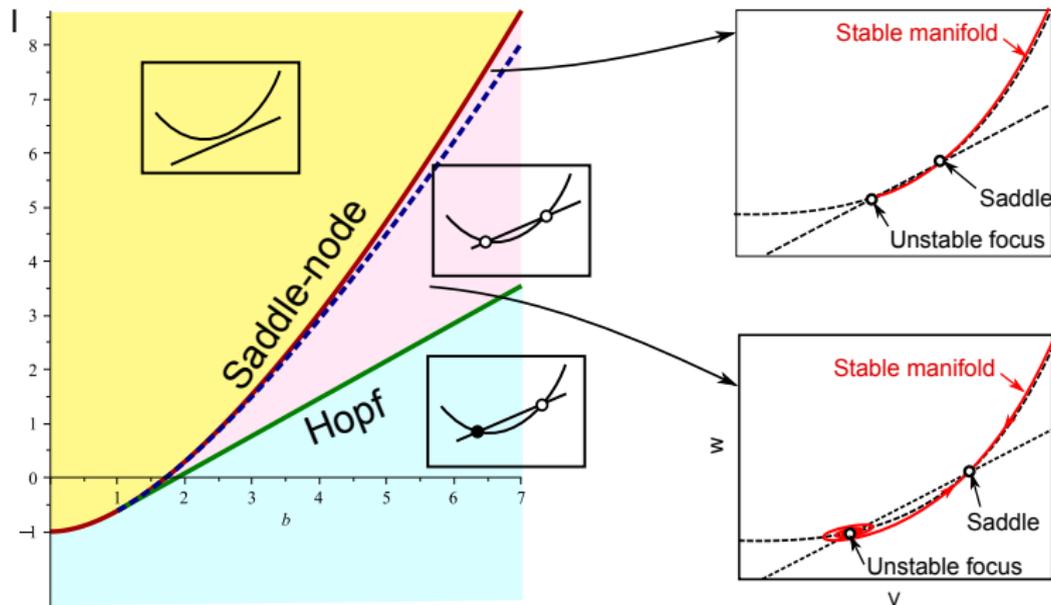
Bromide ion electrode potential in the Belousov–Zhabotinsky reaction; figure from [J.L. Hudson et al., 1979]

Our aim was to show that they also occur in 2D integrate-and-fire models through the simple geometric mechanism.

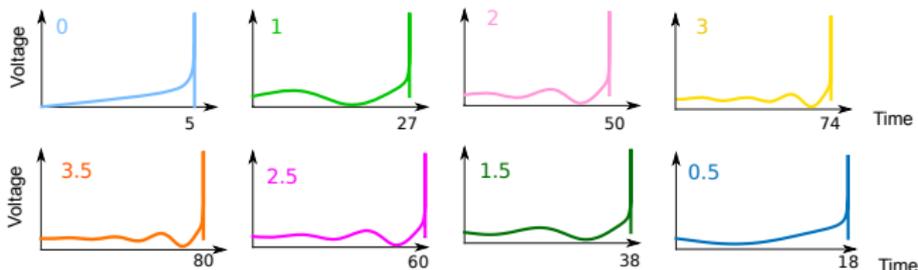
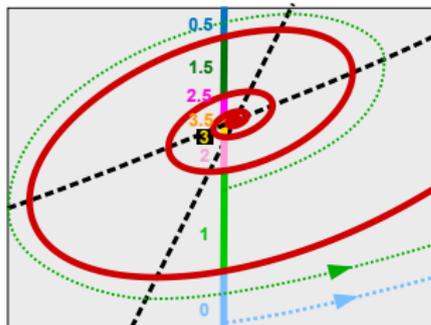
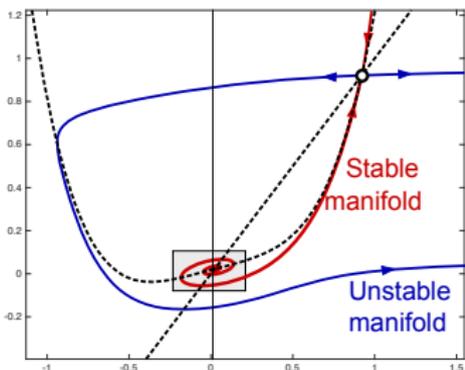
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Mixed-Mode Oscillations (MMOs, slow oscillations interspersed with spikes or bursts):



Bidimensional integrate-and-fire models were studied in [R. Brette, W. Gerstner, 2005], [E. Izhikevich, 2003], [N. Jimenez et al., 2013] and [J. Touboul, R. Brette, 2009].

We assume that:

- $F \in C^3(\mathbb{R})$ (at least)
- F is strictly convex
- $\lim_{v \rightarrow -\infty} F'(v) < 0$
- there exist $\varepsilon > 0$ and $\delta > 0$ such that:

$$\lim_{v \rightarrow \infty} \frac{F(v)}{v^{2+\varepsilon}} \geq \delta$$

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Definition [Adaptation map]

The adaptation map Φ associates to a value of the adaptation variable w the value of the adaptation variable after reset:

$$\Phi(w) := \gamma W(t^*; v_r, w) + d,$$

where $(V(t; v_r, w), W(t; v_r, w))$ is the solution of the system (1)-(2) with initial condition (v_r, w) at time t , and t^* is the value at which $V(t; v_r, w)$ diverges.

Let $\mathcal{D} = \{w_1, w_2, \dots\}$ be the set of intersections of the line $v = v_r$ with SMSFP. Then $\Phi : \mathbb{R} \setminus \mathcal{D} \rightarrow \mathbb{R}$ is well-defined.

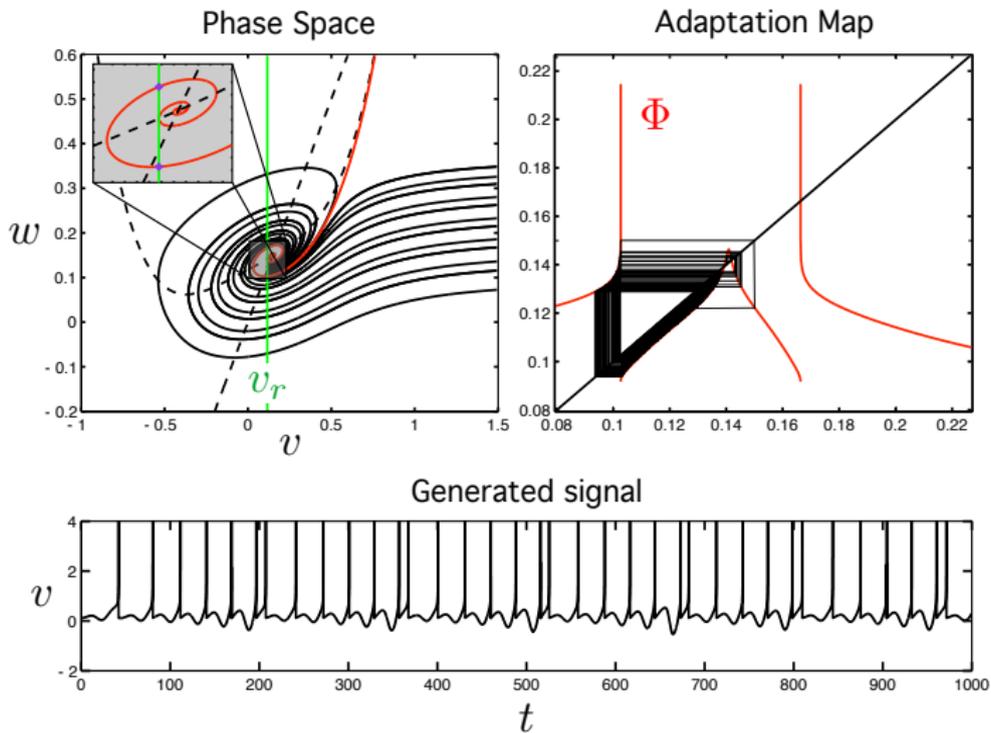
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Remark

Suppose that $(v(t), w(t))$ is the spiking solution starting from (v_r, w_0) and let $\{t_n\}_{n>0}$ be the sequence of spike times for this solution. By $\{w_n\}_{n>0}$ denote the values of the adaptation variable w at spikes, i.e.

$$w_n := w(t_n^+) = \gamma w(t_n^-) + d$$

Then the adaptation map satisfies

$$\Phi(w_n) = w_{n+1}$$

The spike train can be qualitatively described via iterations of Φ , with fixed points of Φ corresponding to tonic, regular spiking and periodic orbits to bursts. Thus the study of the dynamics of Φ allows to discriminate between different spiking patterns.

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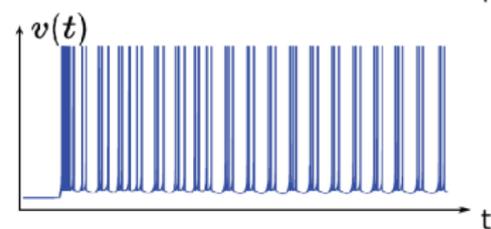
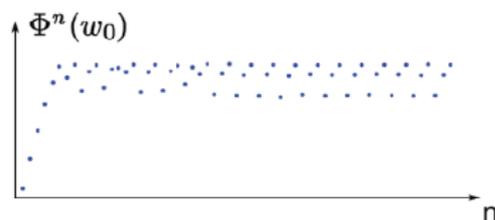
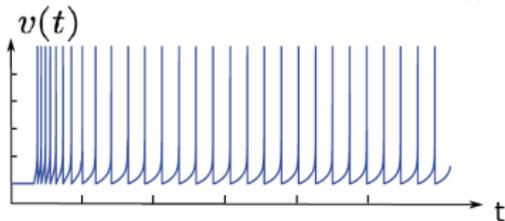
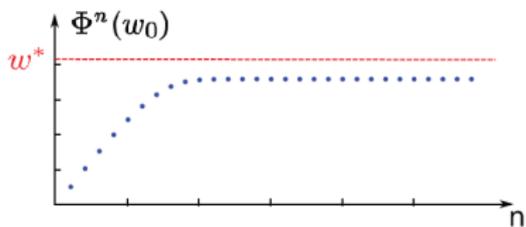
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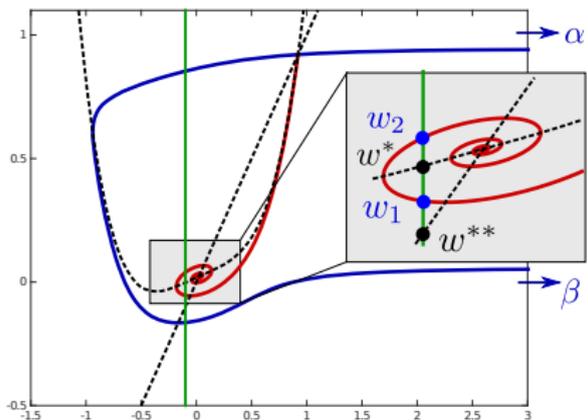
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(figures from [J. Touboul, R.Brette, 2009])



- $(w_i)_{i=1 \dots p}$ - intersections of the reset line $\{v = v_r\}$ with SMSFP
- p_1 - the index such that $(w_i)_{i \leq p_1}$ are below the v -nullcline and $(w_i)_{i > p_1}$ are above
- $(I_j)_{j=0 \dots p+1}$ - intervals with endpoints w_i
- α, β - the value of w after a spike for an initial condition on the upper and, respectively, lower branch of UMSFP

Theorem

The adaptation map has the following properties:

- i it is defined for all $w \in \mathcal{D} = \mathbb{R} \setminus \{w_i; i = 1 \dots p\}$
- ii its regular (at least C^1) everywhere except the points $(w_i)_{i=1 \dots p}$
- iii at the boundaries of the definition domain \mathcal{D} , $\{w_i; i = 1 \dots p\}$, the map has well-defined and distinct left and right limits:

$$\begin{cases} \lim_{w \rightarrow w_i^-} \Phi(w) = \alpha, & \lim_{w \rightarrow w_i^+} \Phi(w) = \beta, & i \leq p_1 \\ \lim_{w \rightarrow w_j^-} \Phi(w) = \beta, & \lim_{w \rightarrow w_j^+} \Phi(w) = \alpha, & j > p_1 \end{cases}$$

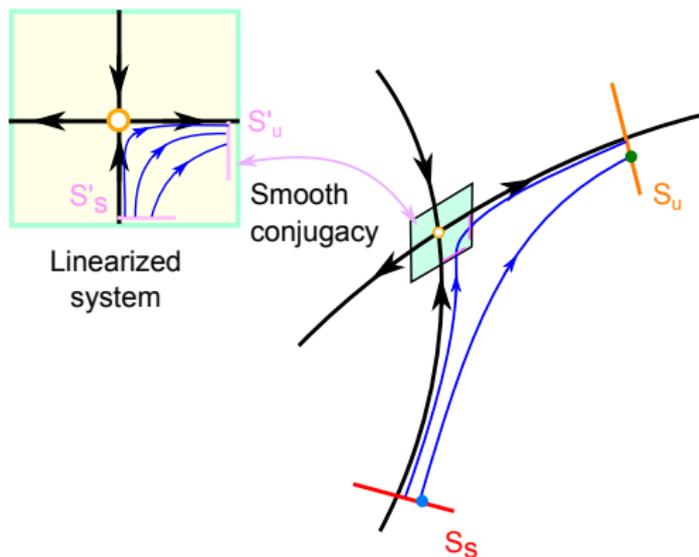
- iv the derivative $\Phi'(w)$ diverges at the discontinuity points:

$$\begin{cases} \lim_{w \rightarrow w_i^\pm} \Phi'(w) = \infty & i \leq p_1 \\ \lim_{w \rightarrow w_i^\pm} \Phi'(w) = -\infty & i > p_1 \end{cases}$$

- Ⓧ for $w < \min\{\frac{d}{1-\gamma}, w_1, w^{**}\}$ we have $\Phi(w) \geq \gamma w + d > w$
- Ⓨ $\Phi(w)$ is convex in the left-neighbourhood of w_i (and concave in the right-neighbourhood)
- Ⓩ $\Phi(w)$ has a horizontal plateau for $w \rightarrow \infty$ provided that

$$\lim_{v \rightarrow -\infty} F'(v) < -a(b + \sqrt{2})$$

The divergence of the derivative $\lim_{w \rightarrow w_1} \Phi'(w) = \infty$ is due to the magnitudes of the eigenvalues $\nu > 0$ and $\mu < 0$ of the saddle fixed point: $|\nu| - \mu > 0$



Assume that the line $v = v_r$ has two intersections with SMSFP: w_1 and w_2 , with $w_1 < w_2$. We distinguish the following cases:

$$I. \beta < w_1 < \alpha < w_2$$

$$I.' \beta < w_1 < w_* < w_2 < \alpha$$

$$II. \alpha < w_* < w_2$$

$$III. \Phi(\beta) \geq \beta$$

$$I.'' \beta < \alpha < w_1$$

$$II.' w_* \leq \alpha < w_2$$

$$III.' \Phi(\beta) < \beta$$

$$I.''' w_1 < \beta < \alpha$$

$$IV.a \Phi(\alpha) \leq \Phi(\beta)$$

$$V.a w_1 < w_2 < \beta < \alpha$$

$$IV.b \Phi(\alpha) > \Phi(\beta)$$

$$V.b w_1 < \beta < w_2 < \alpha$$

$$V.c w_1 < \beta < \alpha < w_2$$

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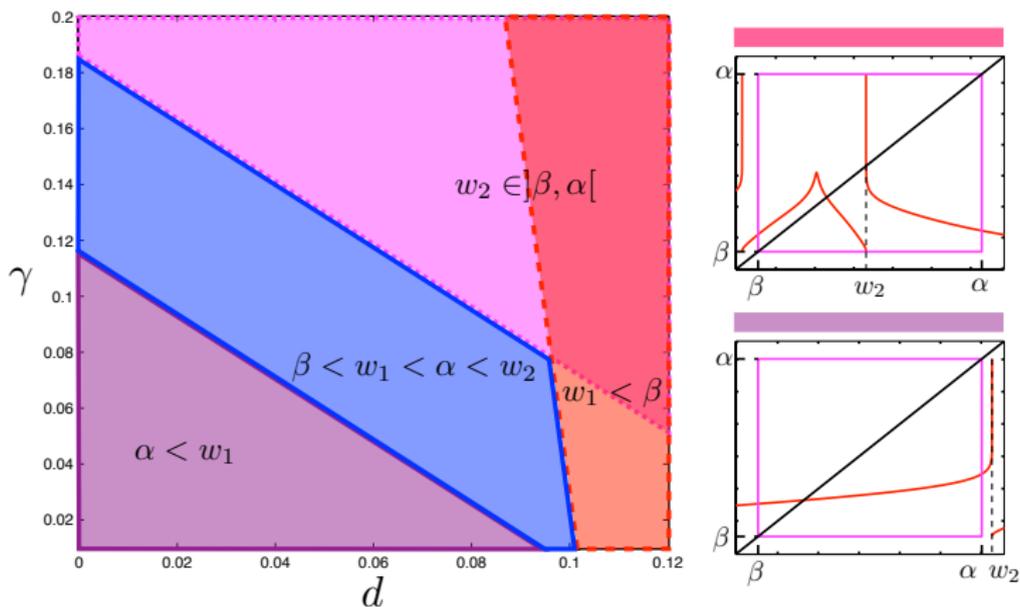
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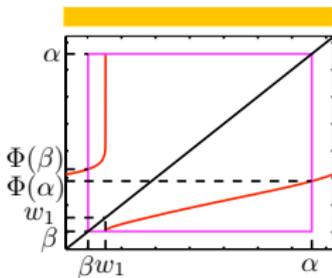
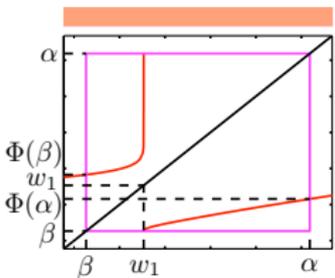
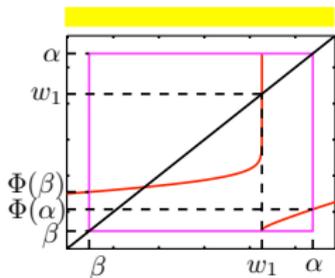
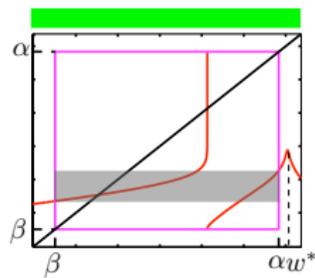
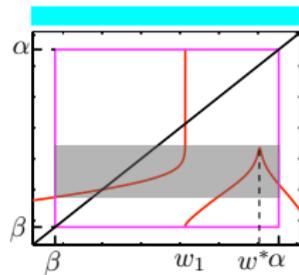
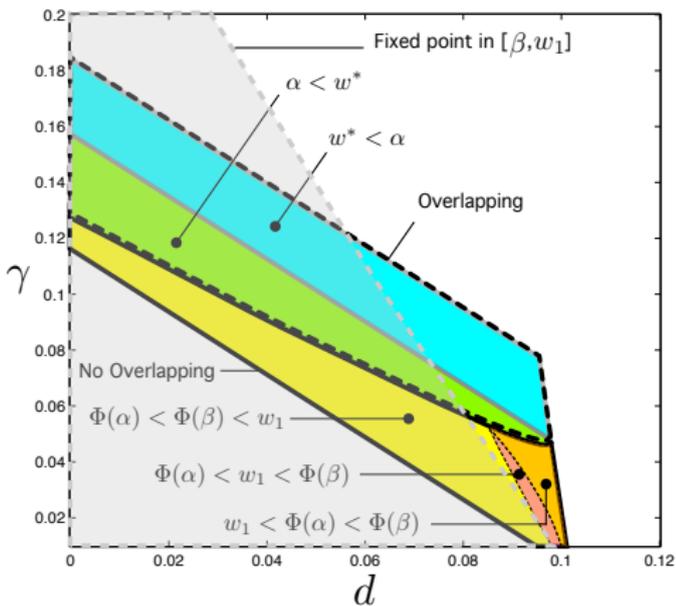
$$\text{V.c } w_1 < \beta < \alpha < w_2$$

Quartic model $F(v) = v^4 + 2cv$ with parameter values: $a = 0.1$, $b = 1$,
 $c = 0.1$, $I = -3(a/4)^{(4/3)}(2a - 1) + 0.1 \approx 0.1175$ and $v_r = 0.1158$



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Firstly, let us assume that

1. $\beta < w_1 < \alpha < w_2$

The above means that the identity line passes through the gap at w_1 and that in the interval $(-\infty, \alpha]$ (where the dynamics concentrates) there is only one discontinuity point w_1 .

The analysis of Φ will be cut to the invariant interval $[x, \alpha]$, where $x = \beta$ when $\Phi(\beta) \geq \beta$ or $x = w_f$ when $\Phi(\beta) < \beta$ and $w_f < \beta$ is the greatest fixed point of Φ in $(-\infty, \beta)$

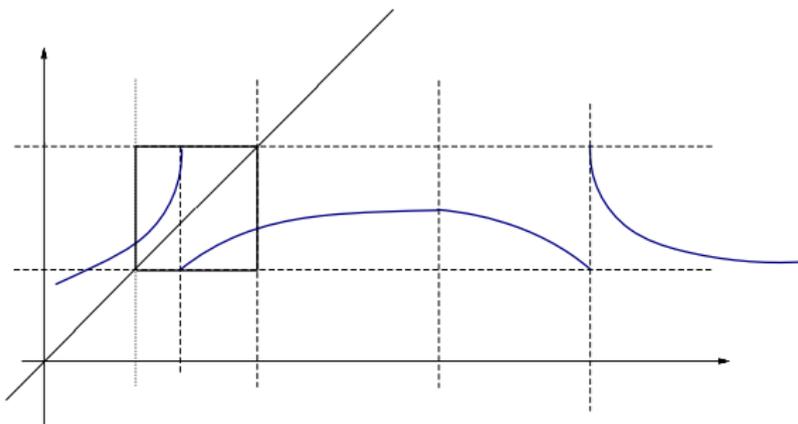
Now let us add the following two assumptions about Φ :

- II. $\alpha < w_* < w_2$
- III. $\Phi(\beta) \geq \beta$

Proposition

Under I. and II., whenever $\Phi : [x, \alpha] \setminus \{w_1\} \rightarrow [x, \alpha]$ has a periodic orbit (with period $q > 1$), this periodic orbit exhibits MMBO. However, this orbit does not need to be stable.

Assume I., II. and III.



We analyze $\Phi : [\beta, \alpha] \rightarrow [\beta, \alpha]$:

- $\Phi(w)$ is piecewise C^1 on $[\beta, \alpha]$ with a single jump discontinuity at $w = w_1 \in (\beta, \alpha)$.
- $\lim_{w \rightarrow w_1^-} \Phi'(w) = \lim_{w \rightarrow w_1^+} \Phi'(w) = \infty$
- $\lim_{w \rightarrow w_1^+} \Phi(w) = \beta$ and $\lim_{w \rightarrow w_1^-} \Phi(w) = \alpha$

According to [J.P.Keener, 1980] analysis of such maps can be performed separately for the following cases:

i) non-overlapping case:

$$\Phi(\alpha) < \Phi(\beta)$$

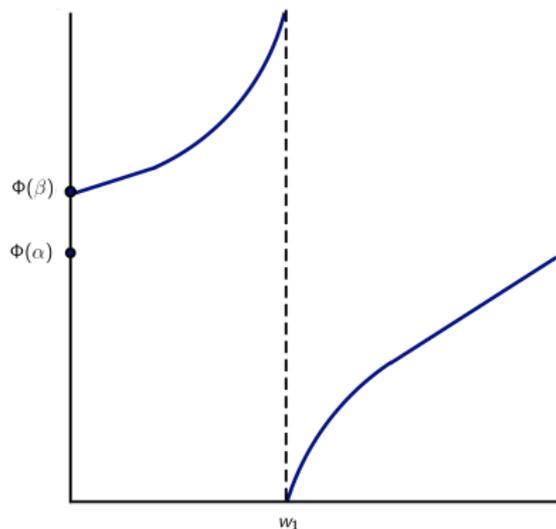
ii) overlapping case:

$$\Phi(\alpha) > \Phi(\beta)$$

iii) $\Phi(\alpha) = \Phi(\beta)$

with the help of the *rotation number*:

$$\varrho(w) := \lim_{n \rightarrow \infty} \frac{\Psi^n(w) - w}{n(\alpha - \beta)} \equiv \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=1}^n \chi_{(w_1, \alpha]}(\Phi^i(w))$$



Non-overlapping case: I., II., III. and IV. a $\Phi(\alpha) < \Phi(\beta)$

We consider the lift $\Psi : \mathbb{R} \rightarrow \mathbb{R}$ of Φ by identifying α with β and requiring that $\Psi(w + k(\alpha - \beta)) = \Psi(w) + k(\alpha - \beta)$, for all $k \in \mathbb{Z}$ and $w \in \mathbb{R}$.

Theorem (cf. [R.Brette, 2003] and [F.Rhodes,Ch.Thompson, 1986])

The rotation number

$$\lim_{n \rightarrow \infty} \frac{\Psi^n(w) - w}{n(\alpha - \beta)} = \varrho$$

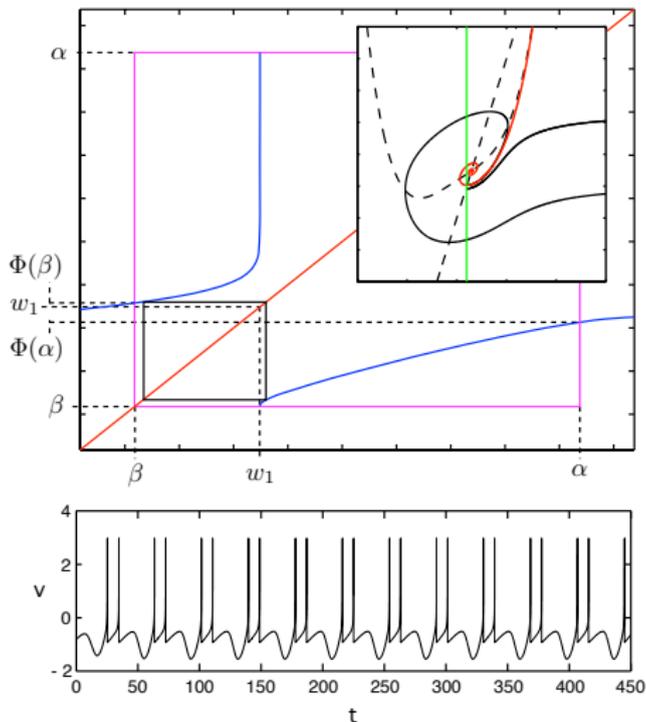
exists and does not depend on $w \in [\alpha, \beta]$.

Moreover, if $\varrho = \frac{p}{q} \in \mathbb{Q}$, then all orbits $\{\Phi^n(w)\}$, $w \in [\beta, \alpha]$, tend to a periodic orbit with the same period q and if $\varrho \notin \mathbb{Q}$, then all orbits have the same limit set which is either the whole $[\beta, \alpha]$ or some Cantor subset of it (meaning, in particular there are no periodic orbits).

- $\varrho = 0 \pmod{1} \implies$ tonic, regular spiking (for every initial condition $w_0 \in [\beta, \alpha] \setminus \{w_1\}$)
- $\varrho = p/q \in \mathbb{Q} \setminus \mathbb{Z} \implies$ MMBO (with periodicity of interspike-intervals and interspersing oscillations)
- $\varrho \in \mathbb{R} \setminus \mathbb{Q} \implies$ no periodic orbits and we observe *chaos*.

Proposition

Under I., II., III., IV.a., if $\Phi(\beta) > w_1$ and $\Phi(\alpha) < w_1$ then Φ has a periodic orbit of period two, which exhibits MMBO.



If $\Psi : \mathbb{R} \rightarrow \mathbb{R}$ is a non-decreasing map of degree-one (i.e. in our case $\Psi(w + (\alpha - \beta)) = \Psi(w) + (\alpha - \beta)$ for every $w \in \mathbb{R}$), then

$$R(\Psi) := \{(x, y) \in \mathbb{R}^2 : \Psi^-(x) \leq y \leq \Psi^+(x)\}.$$

Definition [*H*-convergence]

$\Psi_s \xrightarrow{H} \Psi_{s_0}$ as $s \rightarrow s_0$ iff $R(\Psi_s) \xrightarrow{R} R(\Psi_{s_0})$ in the Hausdorff metric

We say that (Ψ_s) is uniformly convergent to Ψ_{s_0} at x_0 as $s \rightarrow s_0$ if for each $\varepsilon > 0$ there exist $\xi > 0$ and $\delta > 0$ such that for all s and x satisfying $|s - s_0| < \xi$ and $|x - x_0| < \delta$ we have $|\Psi_s(x) - \Psi_{s_0}(x_0)| < \varepsilon$. The *H*-convergence for non-decreasing degree-one circle maps can be characterised in a very convenient way:

Proposition (cf. [F.Rhodes,Ch.Thompson, 1991])

If (Ψ_s) is a family of degree-one non-decreasing maps, then $\Psi_s \xrightarrow{H} \Psi_{s_0}$ as $s \rightarrow s_0$ if and only if (Ψ_s) is uniformly convergent to Ψ_{s_0} at each point of continuity of Ψ_{s_0} .

If $\Psi : \mathbb{R} \rightarrow \mathbb{R}$ is a non-decreasing map of degree-one (i.e. in our case $\Psi(w + (\alpha - \beta)) = \Psi(w) + (\alpha - \beta)$ for every $w \in \mathbb{R}$), then

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Theorem (cf. [R.Brette, 2003], [F.Rhodes,Ch.Thompson, 1991])

Suppose that $s \mapsto \Phi_s$, $s \in [c, d]$, is a family of adaptation maps with strictly increasing lifts Ψ_s such that the mapping $(s, w) \mapsto \Psi_s(w)$ is increasing with respect to each variable and $s \mapsto \Psi_s$ is continuous with respect to the topology of H -convergence. Let ρ_s be the rotation number of Ψ_s . Then:

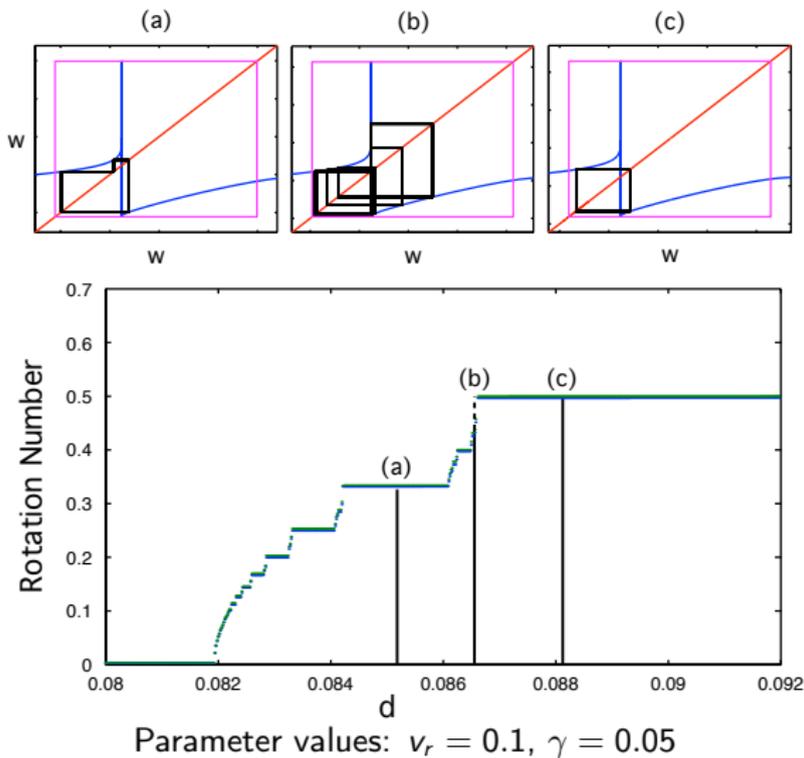
- $\rho : s \mapsto \rho_s$ is continuous and non-decreasing;
- for all $p/q \in \mathbb{Q} \cap \text{Im}(\rho)$, $\rho^{-1}(p/q)$ is an interval containing more than one point, unless it is $\{c\}$ or $\{d\}$;
- ρ reaches every irrational number at most once;
- ρ takes irrational values on a Cantor-type subset of $[c, d]$, up to a countable number of points

Proposition

Let a, b, v_R, I and γ be fixed and consider d varying in some interval $d \in [\lambda_1, \lambda_2]$. Suppose that for this choice of parameter values a, b, v_R, I and γ the adaptation map Φ_d satisfies conditions I., II., III. and IV.a for any value of $d \in [\lambda_1, \lambda_2]$

Let ϱ_d denote the unique rotation number obtained for the map Φ_d (considered on the "fundamental interval" $[\beta_d, \alpha_d]$). Then the mapping $\rho : d \mapsto \varrho_d$ is continuous.

If moreover, for every $d \in [\lambda_1, \lambda_2]$, the adaptation map Φ_d satisfies $\Phi_d(\beta_{\lambda_1}) > \Phi_d(\alpha_{\lambda_2})$, then the above mapping $\rho : d \mapsto \varrho_d$ behaves like a Devil's staircase.



➡ The rotation number $\varrho = p/q \in \mathbb{Q}$ characterises the signature of MM(B)O:

$$\mathcal{L}_1^{s_1}, \mathcal{L}_2^{s_2}, \mathcal{L}_3^{s_3}, \dots$$

where \mathcal{L}_i denotes the number of big oscillations (spikes) and s_i is the number of following them small su threshold oscillations.

For example, $\varrho = 1/3$ corresponds to the periodic signature 3^1 and $\varrho = 3/5$ to the periodic signature $2^1, 1^1, 2^1$

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Overlapping case: I., II. and IV. b $\Phi(\beta) < \Phi(\alpha)$

The analysis of $\Phi : [\beta, \alpha] \rightarrow [\beta, \alpha]$ in the overlapping regime can be made via the results of [M. Misiurewicz, 1986] on *old heavy maps*.

Let $\Psi : \mathbb{R} \rightarrow \mathbb{R}$ denote the lift of $\Phi \upharpoonright [\beta, \alpha]$. The map Ψ is a degree one map with only negative jumps. We can define the following:

Definition [Rotation interval $[a(\Psi), b(\Psi)]$]

$$a(\Psi) := \inf_{w \in \mathbb{R}} \liminf_{n \rightarrow \infty} \frac{\Psi^n(w) - w}{n(\alpha - \beta)},$$

$$b(\Psi) := \sup_{w \in \mathbb{R}} \limsup_{n \rightarrow \infty} \frac{\Psi^n(w) - w}{n(\alpha - \beta)}.$$

An old heavy map does not need to be monotonous in its intervals of continuity and therefore:

Remark

If we assume IV.b, then we can skip the assumption II. since the induced lift Ψ remains an old heavy map.

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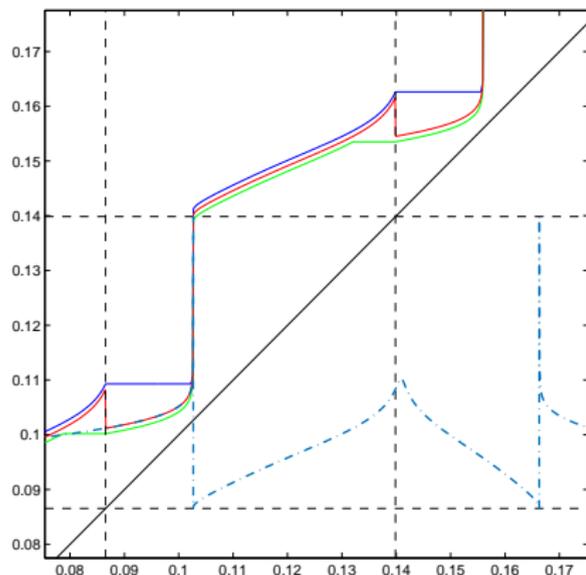
Remark

If we assume IV.b, then we can skip the assumption II. since the induced lift Ψ remains an old heavy map.

Define:

$$\Psi_l(w) := \inf\{\Psi(z) : z \geq w\}$$

$$\Psi_r(w) := \sup\{\Psi(z) : z \leq w\}$$



The maps $\Psi_l(w)$ and $\Psi_r(w)$ are continuous and non-decreasing thus they admit unique rotation numbers.

The nontrivial rotation interval corresponds to complex dynamics (cf. [M. Misiurewicz, 1986]):

- a) if Φ has a q -periodic point w with the rotation number $\varrho(\Psi, w) = p/q$, then $a(\Psi) \leq p/q \leq b(\Psi)$;
- b) if $a(\Psi) < p/q < b(\Psi)$, then Φ has a periodic point w of period q and the rotation number $\varrho(\Psi, w) = p/q$

The coexistence of periodic orbits with infinitely many different periods (non-triviality of the rotation interval) is also sometimes called *chaos* (see [J.P.Keener, 1980]).

Remark

If we additionally assume II., i.e. the monotonicity of Φ in the continuity intervals $[\beta, w_1)$ and $(w_1, \alpha]$, then every such a periodic orbit exhibits MMBO (with both one and no small oscillations between consecutive spikes).

Theorem (cf. [M. Misiurewicz, 1986])

Suppose that the adaptation map $\Phi : [\beta, \alpha]$ is in the overlapping case and that for some ϱ_1 and ϱ_2 we have $a(\Psi) \leq \varrho_1 \leq \varrho_2 \leq b(\Psi)$. Then there exists w_0 such that

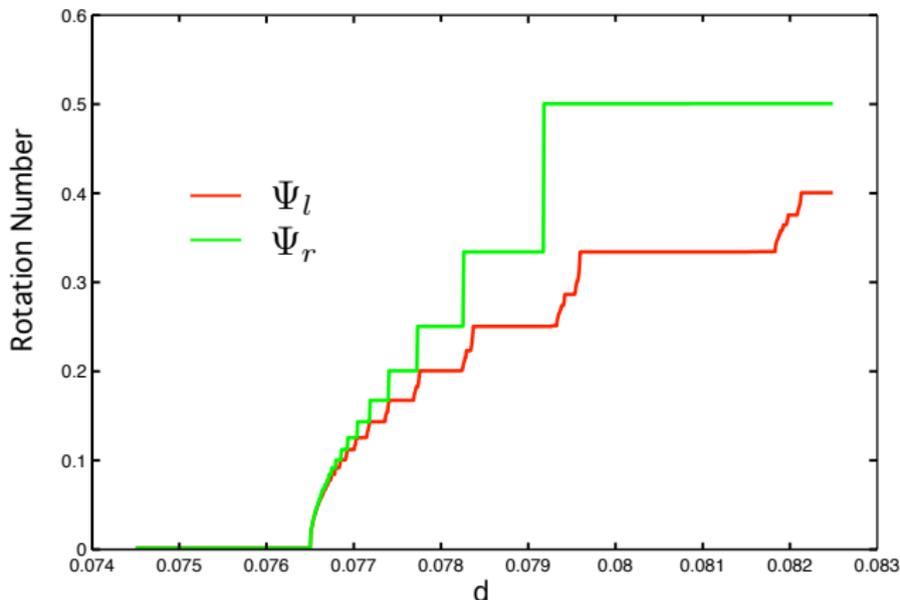
$$\liminf_{n \rightarrow \infty} \frac{\Psi^n(w_0) - w_0}{n(\alpha - \beta)} := \varrho_1$$

$$\limsup_{n \rightarrow \infty} \frac{\Psi^n(w_0) - w_0}{n(\alpha - \beta)} := \varrho_2$$

Proposition

Choose the fixed parameters v_R , a , b , γ and l and the parameter $d \in [\lambda_1, \lambda_2]$ such that for each $d \in [\lambda_1, \lambda_2]$ the map Φ_d is in the overlapping case, i.e. satisfies I., III. and IV.b. Then the maps $d \mapsto a(\Psi_d)$ and $d \mapsto b(\Psi_d)$, assigning to d the endpoints of the rotation interval of Φ_d , are continuous.

Moreover, usually the maps $d \mapsto a(\Psi_d)$ and $d \mapsto b(\Psi_d)$ also behave as Devil's staircase:



Parameter values: $v_r = 0.1$, $\gamma = 0.05$

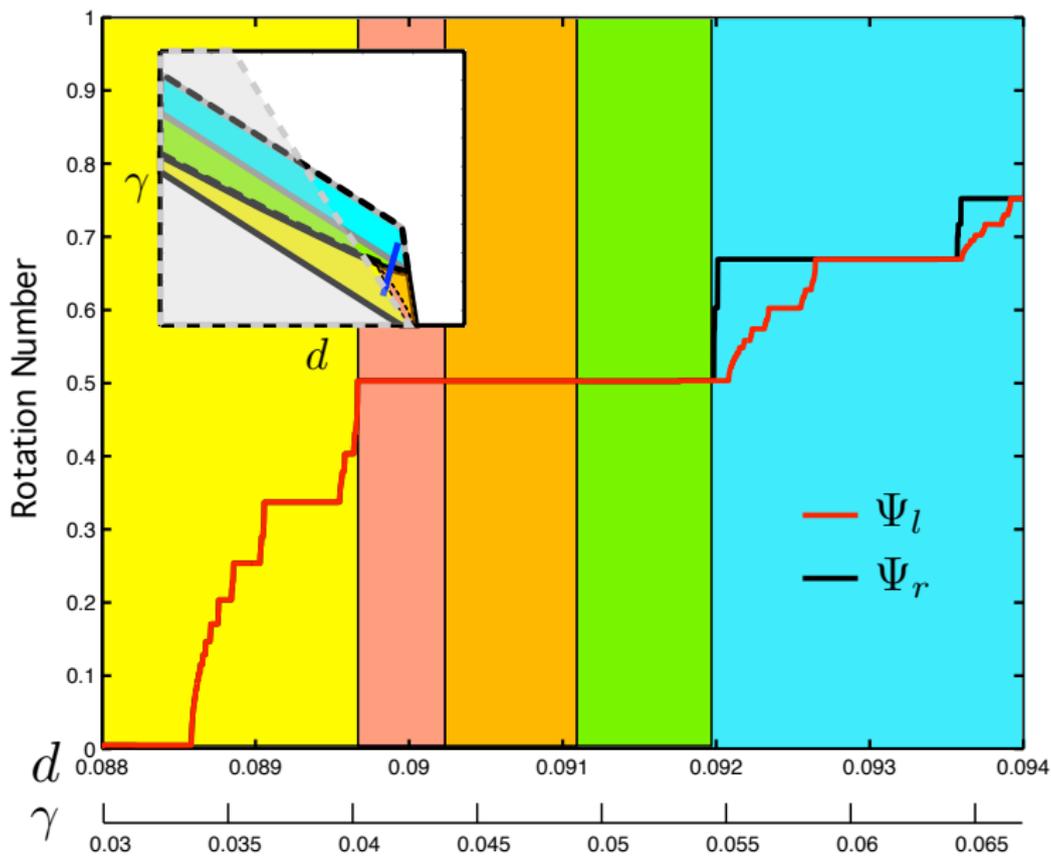
Theorem (Chaos)

*Suppose that Φ satisfies I., II., III and IV. b (an overlapping case with additional monotonicity condition II.). Further assume also that $\Phi(\alpha) < w_1$ and that Φ has at least two periodic orbits, one with period q_1 and the other with period $q_2 \neq q_1$ and that **exactly one** point of each of these periodic orbits is greater than w_1 . Then the mapping $w \mapsto \Phi(w)$ is a **shift on a sequence space**.*

Theorem (Condition for orbits of all periods.)

Existence of a fixed point $w_f \in (\beta, w_1)$ and a periodic orbit with period $q > 1$ implies existence of periodic orbits with arbitrary periods $\tilde{q} > q$ and with MMBO. The same holds if $w_f \in (w_1, \alpha)$ provided that the q -periodic orbit is not of the type q/q (i.e. it admits points to the left and to the right of w_1).

In particular, whenever there is a fixed point $w_f \in (\beta, \alpha)$ and a periodic orbit of the type $1/2$, then there are periodic orbits of all periods, exhibiting MMBO.

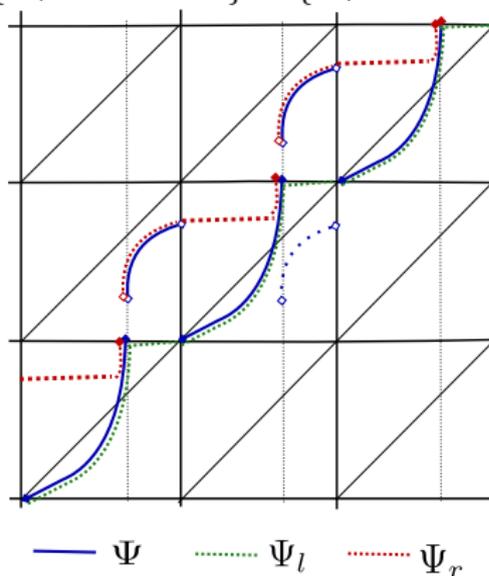


Case of both positive and negative jumps: { I., II. and III.' } or { I., II.' and III.' }

- I. $\beta < w_1 < \alpha < w_2$
- II. $\alpha < w_* < w_2$
- II.' $w_* \leq \alpha < w_2$
- III.' $\Phi(\beta) < \beta$

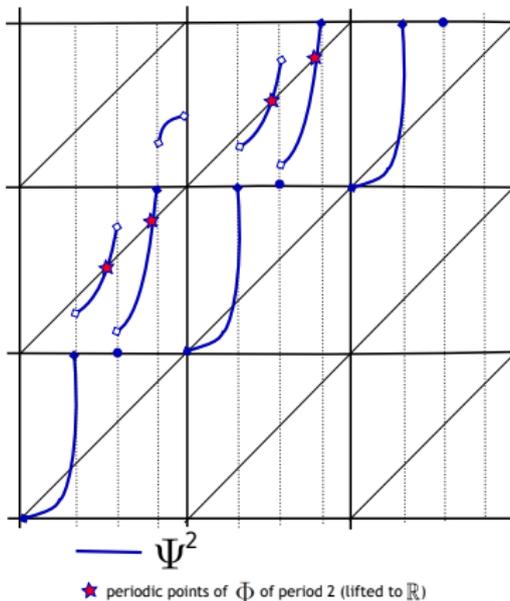
Upper estimate
of the rotation interval:

$$[\varrho(\Psi_l), \varrho(\Psi_r)] \supset [a(\Psi), b(\Psi)]$$



Proposition

Under the assumptions $\{ I., II. \text{ and } III.' \}$ or $\{ I., II.' \text{ and } III.' \}$, if $\Phi(\alpha) > w_1$ and there are no fixed points in (w_1, α) , then Φ has an unstable periodic orbit of period 2. This orbit exhibits MMBO but it is unstable.



No discontinuity points in the invariant interval: I''

Under the following condition

$$I'' \quad \beta < \alpha < w_1$$

there are no discontinuity points w_1 or w_2 in the interval $(-\infty, \alpha)$.

This is the easiest situation:

- since $\Phi(w) > w$ for $w < \min\{\frac{d}{1-\alpha}, w_1, w_{**}\}$ there must be a fixed point in $(-\infty, \alpha)$ and every point w tends under Φ to one of the fixed points. Thus here we observe for every initial condition **tonic, regular spiking** (in particular, we have no MMO and MMBO) and the dynamics is very simple.

No discontinuity points in the invariant interval and the identity line passes below the gap at w_1 : I.'''

$$I.''' \quad w_1 < \beta < \alpha$$

$$V.a \quad w_1 < w_2 < \beta < \alpha$$

Theorem

Suppose that $\lim_{w \rightarrow \infty} \Phi(w) > w_2$. Then every point $w > w_1$ is forward asymptotic either to the fixed point $w_{f,1}$ or to a period two orbit. Under these assumptions no point $w \in \mathbb{R}$ exhibits MMO or MMBO.

No discontinuity points in the invariant interval and the identity line passes below the gap at w_1 : I.'''

$$I.''' \quad w_1 < \beta < \alpha$$

$$V.b \quad w_1 < \beta < w_2 < \alpha$$

Theorem

If $\Phi(w^) < w_2$, then for every $w \in (w_1, w_2)$ we have $\omega(w) \subset \bar{P}$, where \bar{P} denotes the closure of the set of periodic points of $\Phi : [w_1, w_2] \rightarrow [w_1, w_2]$. Particularly, if the set P is finite, then every w tends to some periodic orbit (or fixed point) (with no MM(B)O).*

No discontinuity points in the invariant interval and the identity line passes below the gap at w_1 : I.'''

$$I.''' \quad w_1 < \beta < \alpha$$

$$V.c \quad w_1 < \beta < \alpha < w_2$$

Theorem

Suppose that $\beta < w_ < \alpha$. If*

$$\min \Phi^{-2}(w_*) < \Phi^2(w_*) < \min \Phi^{-1}(w_*) < w_* < \Phi(w_*)$$

and $\Phi(w) > w$ for $w \in (\beta, w_)$, then $\Phi : [\beta, \alpha] \rightarrow [\beta, \alpha]$ has an orbit of period 3. Consequently, Φ has cycles of any period. However, these periodic orbits do not present MMBO.*

Conclusions:

-  We are able to predict the output properties using geometrical analysis
-  In the overlapping and non-overlapping cases existing mathematical tools of rotation theory provide complete description of the dynamics of Φ
-  In the remaining cases (e.g. of both positive and negative jumps) one can obtain weaker results on the dynamics of Φ ; in particular the rotation interval computed via the enveloping maps Ψ_l and Ψ_r gives the upper-estimate for the possible types p/q of periodic orbits

Perspectives:



For multiple discontinuity points the dynamics is even more complex and harder to be completely classified. However, some rigorous results can be obtained via the theory of piece-wise continuous piece-wise monotone maps.



Consider forcing of the IF system through variable I . A simple starting point is a square signal for $I(t)$: the performed analysis can be generalized using a stroboscopic map.



Tackle the problematic of 3D vector field appearing with two recovery variables. In this case we have $\Phi : \mathbb{R}^2 \rightarrow \mathbb{R}^2$. The general mechanism for generating MMBO is the same, yet leading to richer behaviors due to the geometric structure of the flow.

Thank you!

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